

# Soliton-to-Band Optical Absorption in a Quasi One-Dimensional $\text{Pt}^{\text{II}}\text{-Pt}^{\text{IV}}$ Mixed-Valence Complex under Hydrostatic Pressure

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A midgap absorption band is observed in the quasi one-dimensional semiconductor  $[\text{Pt}(\text{en})_2][\text{Pt}(\text{en})_2\text{Cl}_2](\text{ClO}_4)_4$  under hydrostatic pressures at room temperature, where en is ethylenediamine. The transition is allowed only for the polarization parallel to the  $-\text{Cl}-\text{Pt}^{\text{II}}-\text{Cl}-\text{Pt}^{\text{IV}}-$  chain. The peak position remains near the middle of the Peierls gap at any pressure up to 2.2 GPa. The intensity increases exponentially with the peak shift. The gap states responsible for this band are attributed to soliton excitations corresponding to kinks of the charge-density wave.

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The commensurate and nearly commensurate charge-density wave (CDW) systems support a soliton excitation which corresponds to the formation of a domain wall connecting two degenerate phases of periodic distortion of lattice and electron density. Su, Schrieffer, and Heeger have developed a theory of such a kink soliton in long-chain polyenes with the commensurability index of 2.<sup>1</sup> They have evaluated the energy of formation, width, mass, and activation energy for motion of the soliton by the Green's function method taking the lattice discreteness into account explicitly. Subsequently, Takayama, Lin-Liu, and Maki have extended the discrete model to the continuum case which is valid in the limit of weak electron-lattice coupling.<sup>2</sup> These theories show that the soliton state may be described in terms of the order parameter  $\Delta(x) = \Delta_0 \tanh(x/\xi)$ , where  $2\Delta_0$  is the Peierls energy gap and  $\xi$  measures the width of the domain wall. The electronic level of the soliton lies at  $\Delta_0$ , i.e., right at the middle of the Peierls gap. This unique property has been investigated through optical absorption experiments in *trans*-polyacetylene films doped lightly with impurities.<sup>3</sup> Orenstein and Baker have observed the transient optical absorption due to photoexcited solitons in undoped polyacetylene fibrils.<sup>4</sup> This observation has led them to study the recovery kinetics of thermally inequilibrium solitons.<sup>5</sup>

The energy of formation,  $E_s$ , of a soliton is shown to be  $0.6\Delta_0$  in the discrete model,<sup>1</sup> and  $2\Delta_0/\pi \approx 0.64\Delta_0$  in the continuum model.<sup>2</sup> The energy of formation is a fundamental parameter for the characterization of a soliton since it is crucial to statistical properties. Very little experimental verification, however, has been reported on the formation of a kink soliton in an intrinsic material where its physical properties can be controlled by external parameters such as pressure. The significance of ex-

citing a soliton in an intrinsic material lies in the fact that it can be formed anywhere in the medium, and can propagate indefinitely unless it encounters some structural irregularities due to size and/or imperfections such as is the case with *trans*-polyacetylene. In this Letter, we report experimental results of optical absorption which show the presence of solitons in single crystal of  $[\text{Pt}(\text{en})_2][\text{Pt}(\text{en})_2\text{Cl}_2](\text{ClO}_4)_4$ . Here, en refers to the ethylenediamine molecule. A midgap absorption band appears upon the onset of hydrostatic pressure, and its intensity increases markedly with pressure. This new effect is interpreted to arise from an increase in the soliton density due to the decrease in the order parameter, and thus in  $E_s$ .

$[\text{Pt}(\text{en})_2][\text{Pt}(\text{en})_2\text{Cl}_2](\text{ClO}_4)_4$  is one of halogen-bridged mixed-valence platinum complexes (hereafter abbreviated as HMPC). Valence electrons in HMPC form a linear-chain CDW with its commensurability index of 2 in accordance with *trans*-polyacetylene. Pt and Cl ions form either  $-\text{Pt}_n^{\text{II}}-\text{Cl}-\text{Pt}_n^{\text{IV}}-\text{Cl}-\text{Pt}_{n+1}^{\text{II}}-$  or  $-\text{Pt}_n^{\text{IV}}-\text{Cl}-\text{Pt}_n^{\text{II}}-\text{Cl}-\text{Pt}_{n+1}^{\text{IV}}-$  chains which are degenerate in energy. These chains are oriented along the crystallographic *b* axis. They are separated from each other by 8.3 and 9.7 Å towards the (101) and the *c* directions, respectively,<sup>6</sup> so that the interaction between chains is very weak. The electron-lattice interaction within the chain consists of coupling of Pt ions with intervening Cl ions: The charge transfer between neighboring Pt ions is induced by the displacement of the Cl ion from the midpoint of two Pt ions. This is in contrast to the electron-lattice coupling in *trans*-polyacetylene, in which the coupling originates from the mutual displacement of adjacent C atoms. Ichinose has proposed the possibility of the soliton formation in HMPC.<sup>7</sup> Recently, Onodera has examined the physical implication of the

difference in the electron-lattice coupling mechanism between *trans*-polyacetylene and HMPc.<sup>8</sup> According to their theoretical arguments, the Hamiltonian for the halogen-bridged Pt chain becomes equivalent to that for the  $(CH)_x$  chain in *trans*-polyacetylene in the continuum limit. This result assures that the soliton state in HMPc can be described in terms of a model similar to that of *trans*-polyacetylene as far as energetic aspects are concerned. The Peierls gap in  $[Pt(en)_2][Pt(en)_2Cl_2](ClO_4)_4$  is expected to be some 3 eV with reference to the energy 2.8 eV of the charge-transfer exciton associated with this gap.<sup>9,10</sup> Consequently, the soliton state is expected to lie near 1.5 eV above the maximum of the valence band.

The experiment is made by the optical absorption of single crystals subjected to hydrostatic pressure. The single crystal is synthesized by the procedure described elsewhere.<sup>9</sup> The hydrostatic pressure is generated by the use of a diamond-anvil apparatus. The sample is cut into an area of  $50 \times 150 \mu m^2$  and cleaved into a thickness of 10–20  $\mu m$ . It is immersed in the pressure-transmitting fluid in the hole of a metal gasket together with a 10- $\mu m$ -thick bare sheet of Polaroid linear polarizer. The polarizer is pasted by silicon grease on the culet of the upper diamond anvil, so that the direction of polarization with respect to the crystal orientation can be adjusted by rotation of the upper diamond anvil. The absorption spectrum is measured by a microspectroscopy technique with the spot size of 20  $\mu m$  in diameter. The radiation from a microscope is fed to a spectrometer through an optical fiber. This technique allows us to make the sample-in, sample-out measurement without any change in the experimental situation. In particular, the spectrum obtained by this method is not influenced by the depolarization effect due to strain in the diamond at any pressure. All the measurements are performed at room temperature.

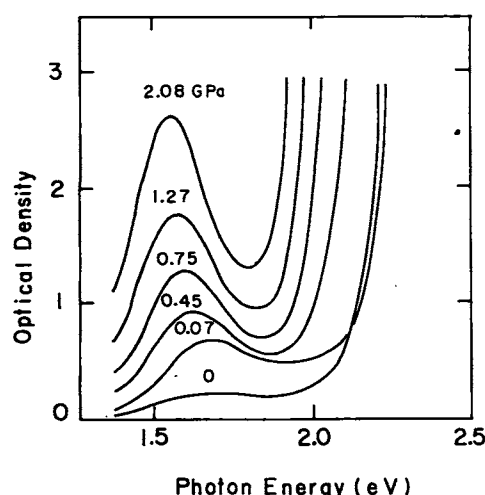


FIG. 1. Pressure dependence of the absorption spectrum of  $[Pt(en)_2][Pt(en)_2Cl_2](ClO_4)_4$  for  $E||b$ .

We measure the spectrum of the optical density, i.e.,  $-\log T$ , for the  $E||b$  polarization, where  $T$  refers to the transmittance and  $E$  to the electric field of radiation. There is a leakage of about 0.8% of the incident radiation throughout any selected area of the sample with no regard to the wavelength. The leakage is independent of pressure. This phenomenon is probably due to the optical inhomogeneity which seems common to many HMPc crystals.<sup>10</sup> Hence we make correction for this leakage to obtain the net spectrum for  $E||b$ . The results obtained for the sample with a thickness of 20  $\mu m$  is shown in Fig. 1.

At atmospheric pressure, the spectrum is dominated by the intense absorption band (CT band) due to the creation of the charge-transfer exciton mentioned above. In addition, a subsidiary structure ( $A$  band) is observed below the CT band. The  $A$  band is intensified by several times upon the onset of a very low pressure ( $\lesssim 0.1$  GPa). The spectral peak lies at the photon energy of 1.67 eV. Then, its intensity increases continuously as the pressure increases. At the same time, it shifts towards lower energies. The absorption edge of the CT band is displaced towards the higher-energy side by 20 meV in coincidence with the discontinuous increase in the intensity of the  $A$  band. Then it shifts continuously towards lower energies with pressure. On release of the pressure, the continuous changes take place reversibly. The  $A$  band remains unremoved, however, even if the pressure is completely released: It is gradually removed if the sample is kept at ambient conditions for a few weeks.

Figure 2 shows the pressure dependence of the shift of the peak position of the  $A$  band and of the photon energy at the optical density of 2.20 of the CT band. Both bands shift nonlinearly with pressure. The magnitude of the shift of the  $A$  band is about 0.4 times that of the CT band in the whole pressure range of the present experiment. It is known that the oscillator strength of the CT

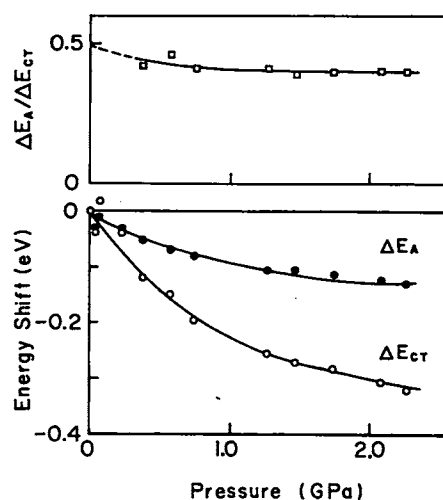


FIG. 2. Pressure dependence of the shift of  $A$  and CT bands.

band increases considerably as the Peierls gap decreases.<sup>10</sup> This effect would raise the apparent rate of the shift of the absorption edge. Consequently, one may safely say that the energy of the *A* band remains near the middle of the Peierls gap regardless of their variations. Moreover, both bands disappear in the  $E \perp b$  polarization. These characteristic behaviors can be explained in terms of the soliton model that the *A* band arises from the optical transition from (to) the occupied (unoccupied) mid-gap level to (from) the conduction (valence) band. Besides, we notice that the spectral line shape is asymmetric with a tail on the higher-energy side if the contribution from the CT band is subtracted. This aspect is also characteristic of the soliton spectrum.<sup>11,12</sup>

The intensity of the *A* band is very sensitive to pressure. The integrated intensity increases rather exponentially as a function of the energy shift as shown in Fig. 3. The oscillator strength amounts to  $1.2 \times 10^{-2}$  at 2 GPa, which corresponds to 0.4% of the oscillator strength, 3.0, of the CT band at the atmospheric pressure.<sup>10</sup> According to theory of the optical absorption in *trans*-polyacetylene, the oscillator strength  $f_s$  of the soliton absorption is approximately given by<sup>11,12</sup>

$$f_s = 2.82(\xi/l)n_s f_0, \quad (1)$$

if the normalized soliton density  $n_s$  is sufficiently low, where  $l$  represents the distance between two adjacent CH groups along the chain direction and  $f_0$  the oscillator strength of the interband absorption in the absence of a soliton. In HMPC,  $l$  corresponds to the distance between adjacent  $\text{Pt}^{\text{II}}$  and  $\text{Pt}^{\text{IV}}$  ions. Applying Eq. (1) to our experimental results, we have  $(\xi/l)n_s \approx 3 \times 10^{-4}$  at atmospheric pressure.

The Peierls gap opens between  $d_{z^2}$  states of  $\text{Pt}^{\text{II}}$  and  $\text{Pt}^{\text{IV}}$  which have bonding and antibonding characters in connection with the  $p_z$  states of Cl, respectively.<sup>9</sup> The dimerization is driven by the collaboration of the electron-electron repulsion between the neighboring Pt ions and the electron-lattice coupling. The kinetic energy and the on-site electron-electron repulsion compete with these interactions. Nasu has developed a mean-field theory of this system, and discussed its optical proper-

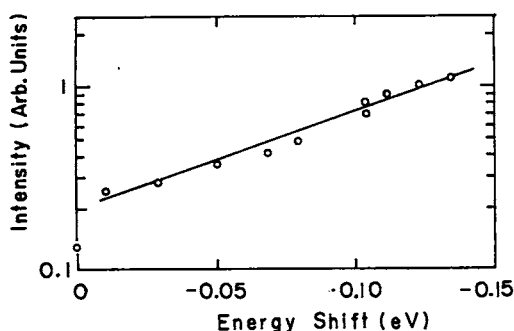


FIG. 3. Plot of the integrated intensity vs the shift of the *A* band.

ties.<sup>13</sup> According to his results, the energy gap decreases quadratically with the transfer energy,  $t_0$ , of the  $d_{z^2}$  orbit between Pt ions. The transfer energy increases upon contraction of the chain, causing the system to be more metallic. Therefore the energy gap and  $\Delta_0$  decrease with pressure. Tanino *et al.* have also observed a pressure-induced red shift of the charge-transfer absorption edge in Wulffram's red salt, which is another member of HMPC.<sup>14</sup>

$\xi/l$  in Eq. (1) can be rewritten in the Takayama-Lin-Liu-Maki model as  $2t_0/\Delta_0$ . The fractional decrease in  $\Delta_0$  can be estimated to be  $\approx 7\%$  at 2 GPa from Fig. 2. A similar amount of increase in  $t_0$  is expected in the first-order approximation. Consequently, the fractional increase in  $\xi/l$  is expected to be (10–20)% at 2 GPa, which is negligibly small in comparison with the increase in  $f_s$ . It follows from these estimates that the change in  $f_s$  induced by pressure arises primarily from the change in the soliton density. The behavior of the integrated intensity of the *A* band (Fig. 3) suggests that the density is limited by an activation energy which decreases as  $\Delta_0$  under pressure. If the activation energy is assumed to correspond to the formation energy, the slope of the experimental points of Fig. 3 yields  $E_s/\Delta_0 = 0.32$ , which is comparable with theoretical values ( $\sim 0.6$ ) mentioned at the beginning of this Letter.

There are two mechanisms of the thermal excitation of solitons. One is the ordinary soliton-antisoliton pair excitation on the uniform part of a chain, where a single soliton cannot be created because of topological restrictions.<sup>11,12</sup> The other is the injection of a single soliton from an edge of the chain. A soliton and its antisoliton can be created individually but alternately at the edge because the topological restriction is lifted by the discontinuity of bonding. Though it is difficult to determine which is the case from our experimental results, one cannot neglect the single-excitation process because the crystal of  $[\text{Pt}(\text{en})_2][\text{Pt}(\text{en})_2\text{Cl}_2](\text{ClO}_4)_4$  consists of columnar structures with a length of 5 to 8  $\mu\text{m}$  and a thickness of about 1  $\mu\text{m}$ .<sup>15</sup> We recall here that the *A* band appears discontinuously at a very low pressure. This fact indicates that the midgap absorption is promoted by a slight modification of lattice. It is likely that the initial compression of the crystal induces a modification of the atomic arrangement near the edge of each columnar structure such as to diminish the potential barrier against the injection of a soliton. The *A* band is occasionally observed even for virgin samples. It does not disappear if the sample is cooled down to liquid-nitrogen temperatures.<sup>9</sup> This fact implies that solitons are confined within chains. This is plausible because the residual potential barrier and/or the activation energy for motion of a soliton could be significantly large in comparison with thermal energies at low temperatures.

In conclusion, the optical absorption band which appears within the Peierls gap in  $[\text{Pt}(\text{en})_2][\text{Pt}(\text{en})_2\text{Cl}_2]$ -

(ClO<sub>4</sub>)<sub>4</sub> exhibits a unique behavior under hydrostatic pressure. The peak position, the spectral line shape, the polarization property, and the pressure dependences of the intensity and the peak position are consistently explained in terms of the soliton model which is similar to the case of *trans*-polyacetylene. The pressure dependence of the intensity suggests that the kink solitons are created predominantly at the edge of the halogen-bridged Pt chain. The detail of the excitation mechanism is not, however, clear yet because we have very little information on crystallographic properties involving the lattice imperfections in this material. A firm identification must, therefore, await experimental studies of the microscopic structure under pressure.

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